Effects of Chain Length on the Structure and Dynamics of Polymer-NP Composites

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Introduction

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The addition of nanoparticles (NP) as fillers to glassforming polymer materials can substantially change the mechanical, dynamical, electrical, and optical properties of the material, thus understanding polymer-NP composite is important for developing new materials. Previous work from the Starr lab has shown that the changes to the glass transition in the dilute NP limit arise from modifications to interfacial dynamics of monomers². In many the applications of these materials, the case is more complex due to higher NP concentrations and longer chain lengths relative to the NP separation. The dilute and concentrated limits have been thoroughly studied, and it has been proposed that there is an intermediate regime, in which NPs can affect each other through chains that bridge between them¹. Therefore, using computational approach, we focus on the intermediate regime, aiming to investigate the chain length dependence of its properties as well as the potential bridging phenomenon.

Chain Structure and Orientation



We first examine the gyration tensor to characterize the chain dimension. Fig. 3a shows that the radius of gyration R_{a^2} increases near the NP surface, meaning that chains elongate approaching the NP. When we scale R_{a^2} by the corresponding average for each chain length $\langle R_{a}^{2} \rangle$ (Fig. 3b), we show that this effect is independent of chain length. We also evaluate the orientational parameters $<P_2\cos\theta_1>$ and $<P_2\cos\theta_3>$ (Fig.4), which are associated with the shortest and longest rotational axis of the chain. Near the NP surface, $\langle P_2 \cos \theta_1 \rangle$ increases while $<P_2\cos\theta_3>$ decreases, demonstrating that chains tend to align their shortest axis normal to the NP surface, and the longest axis along the NP surface. Also, the orientational changes extends farther from the NP as chain length increases.



Simulation Setup

When differentiating bridging chains from others, we find that bridging chains extend along the direction connecting NPs (Fig. 6). However, the differences are only significant when NP separation is large and bridging chains have to stretch, in which cases the bridging chains are rare (see Fig. 5 for the fraction of bridging chains).



We use LAMMPS Molecular Dynamics Simulator to model the polymer composites as an ideal, uniform dispersion of NP surrounded by polymer melts. An icosahedron shaped NP is fixed at the center of a simulation box and surrounded by polymer chains of varying lengths (5, 10, 15, 20, 30, 40 beads); by periodic boundary condition, this simulates an ideal dispersion of NP. For each chain length, we simulate three NP concentrations (2.8%, 5.5%, 10.4%) by varying the total number of polymer chains. For systems with longer chains (N≥30 in ϕ = 2.8%, 5.5%, and N≥20 in ϕ = 10.4%), we use larger simulation cells with 8 fixed NPs surrounded by 8 times the number of polymer chains that were present in the 1 NP systems. These larger simulations ensure that the simulation cell size is always more than double the chain dimension, and allow observation of bridging chains between different NPs. Simulations are performed in an NVT ensemble along an isobaric path with pressure P = 0.1 and temperatures ranging from 0.42 to 0.80, in reduced Lennard-Jones units $T^* = K_b T/\epsilon$, where ϵ is the polymerpolymer interaction strength.





Chain Dynamics and T_q

To evaluate the effect of bridging on the dynamics, we compare the relaxation time τ of bridging chains with that of other chains (Fig.7). Bridging chains have a longer relaxation time compared to the average over all chains and all other subtypes. However, as chain length and the bridging fraction increase, the difference between τ of bridging and that of the average decreases. In other words, the dynamics of bridging chains only differ significantly when bridging between NP requires that chains are significantly expanded. In this case, there are relatively few bridging chains, thus their effect on the average is minimal.

Fig.7 Relaxation time τ (N) for different types of chains.

We then evaluate how chain length and bridging influence the dynamical glass transition temperature T_{q} of composites, based on the composite relaxation time. We show that T_q of composites follows the trend in pure polymer systems (Fig.8). Independent of ϕ , T_q increases with chain length and saturates at N ≈ 20, consistent with previous experiments and simulations of pure polymer systems; we also show that T_a follows a generally linear relationship with 1/N. Since the relaxation of bridging chains only differs from the overall significantly when the fraction of bridging chains is small, the dynamical T_a is not significantly affected by bridging.



Fig.2 Simulation Setup: (a) 1 NP in polymer melts; (b) 8 NPs and bridging chains.

References

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Fig.8 (a) T_g (N) and (b) T_g (1/N) for all different ϕ .

Conclusion

- Bridging chains can have significantly altered conformations and relaxation, but their effects on the composite are minimal.
- Chain length does not have a significant impact on the chain structure at the NP interface, but NP interfacial effects on chain alignment extend farther from the NP surface as chain length increases.
- f_a of composites increases and saturates with chain length, an effect that stems from the chain length effect found in pure polymer systems.